



Thermodynamics of Quantum Information Flows Phys. Rev. Lett. 122, 150603 (2019)

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Quantum Thermodynamics for Young Scientists

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- Topic: relation between the second law of thermodynamics and the information theory
- More specifically: relation between the entropy balance of a subsystem of a Markovian open quantum system and the information flow between the subsystems
- Introduction: second law of thermodynamics for open quantum systems out of equilibrium
- Main result: local Clausius inequality
- Thermodynamics of information flow
- Derivation of the local Clausius inequality
- Application: autonomous quantum Maxwell demon

Second law of thermodynamics for open systems

- Second law of thermodynamics \rightarrow thermodynamics irreversibility
- For open systems out of equilibrium [M. Esposito et al., New J. Phys. 12, 013013 (2010)]

$$\sigma = \Delta S - \sum_{\alpha} \beta_{\alpha} Q_{\alpha} \ge 0$$

where

- σ entropy production
- $S = -\text{Tr}(\rho \ln \rho) \text{von Neumann}$ entropy of the system
- Q_{α} heat delivered from the reservoir α



- Markovian system no memory effects due to correlation with environment – for weak coupling
- Second law in differential form

$$\dot{\sigma} = d_t S - \sum_{lpha} eta_{lpha} \dot{Q}_{lpha} \ge 0$$

where

- $\dot{\sigma}$ entropy production rate
- S = -Tr(ρ ln ρ) von Neumann entropy of the system
- \dot{Q}_{α} heat flow from the reservoir α



Main result: local Clausius inequality

$$\hat{H}_{S} = \sum_{i} \hat{H}_{i} + \hat{H}_{int}$$

- Can we define 2nd law for a single subsystem? Yes, we can!
- Local Clausius inequality

$$\dot{\sigma}_i = d_t S_i - \sum_{\alpha_i} \beta_{\alpha_i} \dot{Q}_{\alpha_i} - \dot{I}_i \ge 0$$

where

- $\dot{\sigma}_i$ local entropy production rate
- S_i = -Tr(ρ_i ln ρ_i) von Neumann entropy of the subsystem i
- Q_{α_i} heat flow from reservoir α_i
- *I_i* information flow between the subsystems (defined later)



Maxwell demons

 Maxwell demon – entropy of a stochastic system can be reduced by a feedback control by an intelligent being



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- Experimental realizations
 - Colloidal particles [S. Toyabe et al., Nat. Phys. 6, 988-992 (2010)]
 - Single-electron boxes [J. V. Koski et al., Phys. Rev. Lett. 113, 030601 (2014)]
 - Superconducting circuits [Y. Masuyama *et al.*, Nat. Commun. 9, 1291 (2018)]
- Generalized second laws [e.g. T. Sagawa, M. Ueda, Phys. Rev. Lett. 100, 080403 (2008)]

$$\Delta S - \Delta I > 0$$

where ΔI – information collected by the feedback device

Autonomous Maxwell demons

- Autonomous Maxwell demon entropy of a stochastic system can be reduced by a feedback control by another stochastic system
- Theoretical proposal Coulomb coupled quantum dots – P. Strasberg *et al.*, Phys. Rev. Lett. 110, 040601 (2013)
- Experimental realization J. V. Koski et al., Phys. Rev. Lett. 115, 260602 (2015)
- Theoretical framework for classical systems based on stochastic thermodynamics – J. M. Horowitz, M. Esposito, Phys. Rev. X 4, 031015 (2014)





Theoretical framework - bipartite dynamics

- Two subsystems: X and Y
- Classical rate equation for state probabilities

$$\dot{p}(x,y) = \sum_{x',y'} \left[W_{x,x'}^{y,y'} p(x',y') - W_{x',x}^{y',y} p(x,y) \right]$$

$$W^{y,y'}_{x,x'}$$
 – rate of transition $(x',y') \rightarrow (x,y)$

 Transitions – either in X or Y, not simultaneous

$$W_{x,x'}^{y,y'} = \begin{cases} w_{x,x'}^y & x \neq x; y = y' \\ w_x^{y,y'} & x = x', y \neq y' \\ 0 & x \neq x', y \neq y' \end{cases}$$



 T_D, μ_D

J. M. Horowitz, M. Esposito, Phys. Rev. X 4, 031015 (2014)

Local second law of thermodynamics

$$I = \sum_{x,y} p(x,y) \ln \frac{p(x,y)}{p(x)p(y)}$$

• Mutual information – measure of correlation between subsystems • Decomposition: $d_t I = \dot{I}_X + \dot{I}_Y$

$$\begin{split} \dot{l}_{X} &= \sum_{x \geq x'; y} \left[w_{x,x'}^{y} p(x',y) - w_{x',x}^{y} p(x,y) \right] \ln \frac{p(y|x)}{p(y|x')} \\ p(x|y) &= p(x,y) / \sum_{y'} p(x,y') \end{split}$$

$$\dot{\sigma}_i = \dot{H}_i - \beta_i \dot{Q}_i - \dot{I}_i \ge 0$$

All already done? Not yet

- Previous approach limited to restricted class of systems
- Only for non-coherent systems (without unitary dynamics)
- Even in classical limit not applicable to systems with $[\hat{H}_{int},\hat{H}_i] \neq 0$
 - Rate equations describe transtions between eigenstates of the total Hamiltonian \hat{H}_{S}
 - For non-commuting Hamiltonians eigenstates of H
 _S are not products of eigenstates of subsystem Hamiltonians H
 _i
 - As a result, rate equations do not have a bipartite structure
- **Our result**: generalizes the concept of autonomous information flow to a generic Markovian open quantum system

Derivation: Assumptions

 Dynamics described by Lindblad equation (internal unitary dynamics + dissipation)

$$d_t \rho = -i \left[\hat{H}^{\mathsf{eff}}, \rho \right] + \mathcal{D}\rho$$

 Additivity of dissipation – interaction with each reservoir gives an independent contribution to the dissipation – valid for Markovian systems

$$\mathcal{D} = \sum \mathcal{D}_{\alpha}$$

Local equilibration

$$\mathcal{D}_{\alpha}\rho_{\alpha}^{\text{eq}} = 0$$
$$\rho_{\alpha}^{\text{eq}} = Z_{\alpha}^{-1} e^{-\beta_{\alpha} \left(\hat{H}_{\text{S}} - \mu_{\alpha}\hat{N}\right)}$$

where $\rho^{\rm eq}_{\alpha}$ – Gibbs state with respect to α

Partial Clausius inequality

• Applying Spohn's inequality [H. Spohn, J. Math. Phys. (N.Y.) 19, 1227 (1978)]

$$-\mathrm{Tr}\left[\left(\mathcal{D}^{lpha}
ho
ight)\left(\ln
ho-\ln
ho_{\mathrm{eq}}^{lpha}
ight)
ight]\geq0$$

one obtains the **partial Clausius inequality** [G. B. Cuetara, M. Esposito, G. Schaller, Entropy 18, 447 (2016)]

$$\dot{\sigma}_{lpha} = \dot{S}^{lpha} - eta_{lpha} \dot{Q}_{lpha} \ge 0$$

where

- $\dot{\sigma}_{\alpha}$ partial entropy production rate
- $\dot{S}^{\alpha} = -\text{Tr}\left[\left(\mathcal{D}^{\alpha}\rho\right)\ln\rho\right]$ rate of change of the von Neumann entropy due to interaction with the reservoir α
- $\dot{Q}_{\alpha} = \text{Tr}\left[\left(\mathcal{D}^{\alpha}\rho\right)\left(\hat{H}_{S}-\mu_{\alpha}\hat{N}\right)\right]$ heat flow from the reservoir α
- **Meaning**: interaction with each reservoir gives a non-negative contribution to the entropy production

Local Clausius inequality

$$\dot{\sigma}_{\alpha} = \dot{S}^{\alpha} - \beta_{\alpha} \dot{Q}_{\alpha} \ge 0$$

 Local entropy production rate – sum of σ
 associated with reservoirs α_i coupled to subsystem i



• We obtain inequality analogous to the result of Horowitz and Esposito

$$\dot{\sigma}_i = d_t S_i - \sum_{lpha_i} eta_{lpha_i} \dot{Q}_{lpha_i} - \dot{I}_i \geq 0$$

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- What is a physical meaning of \dot{I}_i ? Is it related to information flow?
- Answer: yes, it is!

$$\sum_{i}\dot{I}_{i}=d_{t}I$$

where $I = \sum_{i} S_{i} - S$ – (multipartite) mutual information between subsystems

• In the classical limit (secular approximation) with $[\hat{H}_{int}, \hat{H}_i] = 0$ our inequality is equivalent to the result of Horowitz and Esposito

Autonomous quantum Maxwell's demon



K. Ptaszyński, Phys. Rev. E 97, 012116 (2018)

- Operation based on coherent spin exchange + spin selective dissipative dynamics (polarized leads)
- Essentially non-bipartite dynamics: spin exchange simultaneously flip spins in both dots; $[\hat{H}_{int}, \hat{H}_i] \neq 0$
- Could not be described by previously existing approaches

Demon: operation



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Thermodynamics of Quantum Information

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Results

- For J ≤ 100 "pure" Maxwell demon:
 - subsystem 2 cooled $(-\dot{Q}_2 < 0)...$
 - ...with a negligible energy flow $\dot{E}_i \approx 0...$
 - ...due to information flow $(-\dot{Q}_2 > T\dot{I}_2)$ – conversion of heat into work due to feedback control
 - This is compensated by heat dissipation in the first dot $(-Q_1 > T\dot{I_1} > 0)$



- We derived an inequality describing local entropy and information balance of a subsystem of a Markovian open quantum systems coupled to several reservoirs
- This provides a consistent mathematical description of thermodynamics of autonomous information flow in quantum systems
- Applicability of our approach was demonstrated on the example of an autonomous quantum Maxwell demon

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Thank you for your attention!