

Dissipation close to equilibrium

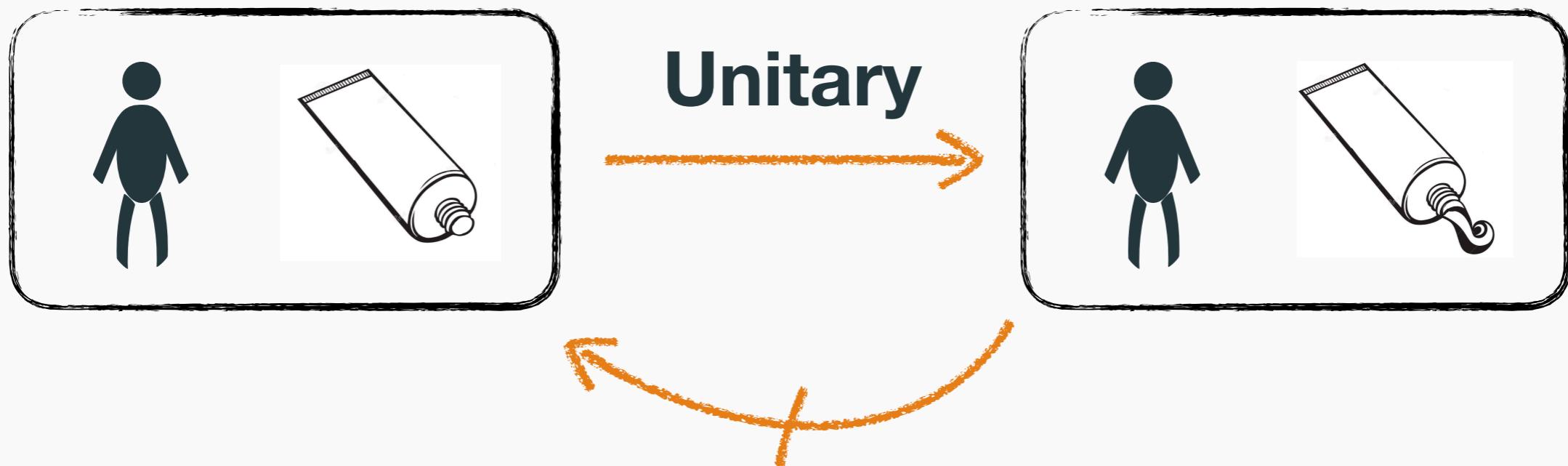
Matteo Scandi



ICFO^R

Irreversibility

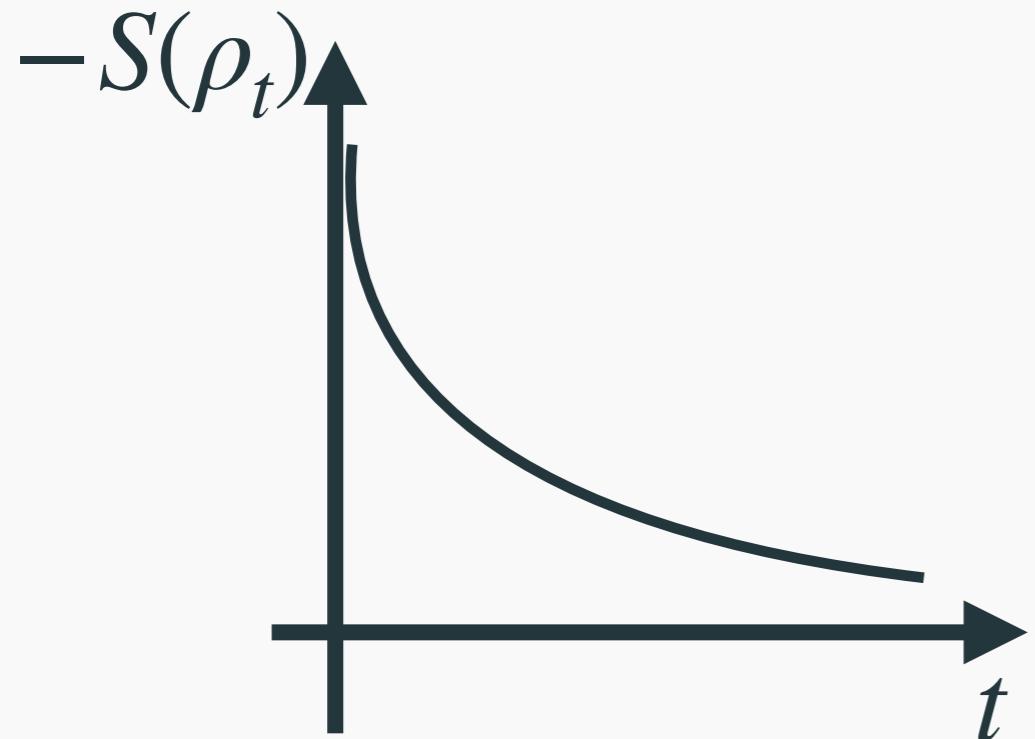
The arrow of time



Entropy: why you can't get the toothpaste back in the tube¹

¹ Whatever works, W. Allen (2009)

Second law of thermodynamics



H theorem: $-S(\rho_t)$ decreases monotonically

Quantum second laws



$\pi :=$ Gibbs state

→ measure of
athermality

H theorem



$S_\lambda(\rho || \pi)$

decreases

$S_\lambda(\rho || \mathcal{D}(\rho))$

decreases

$\mathcal{D}(\rho) :=$ diagonal state

→ measure of
coherence



Framework



$$\rho_S = \text{Tr}_{U \setminus S} [\rho]$$

$$H(t) = H_S(t) + gV + H_B$$

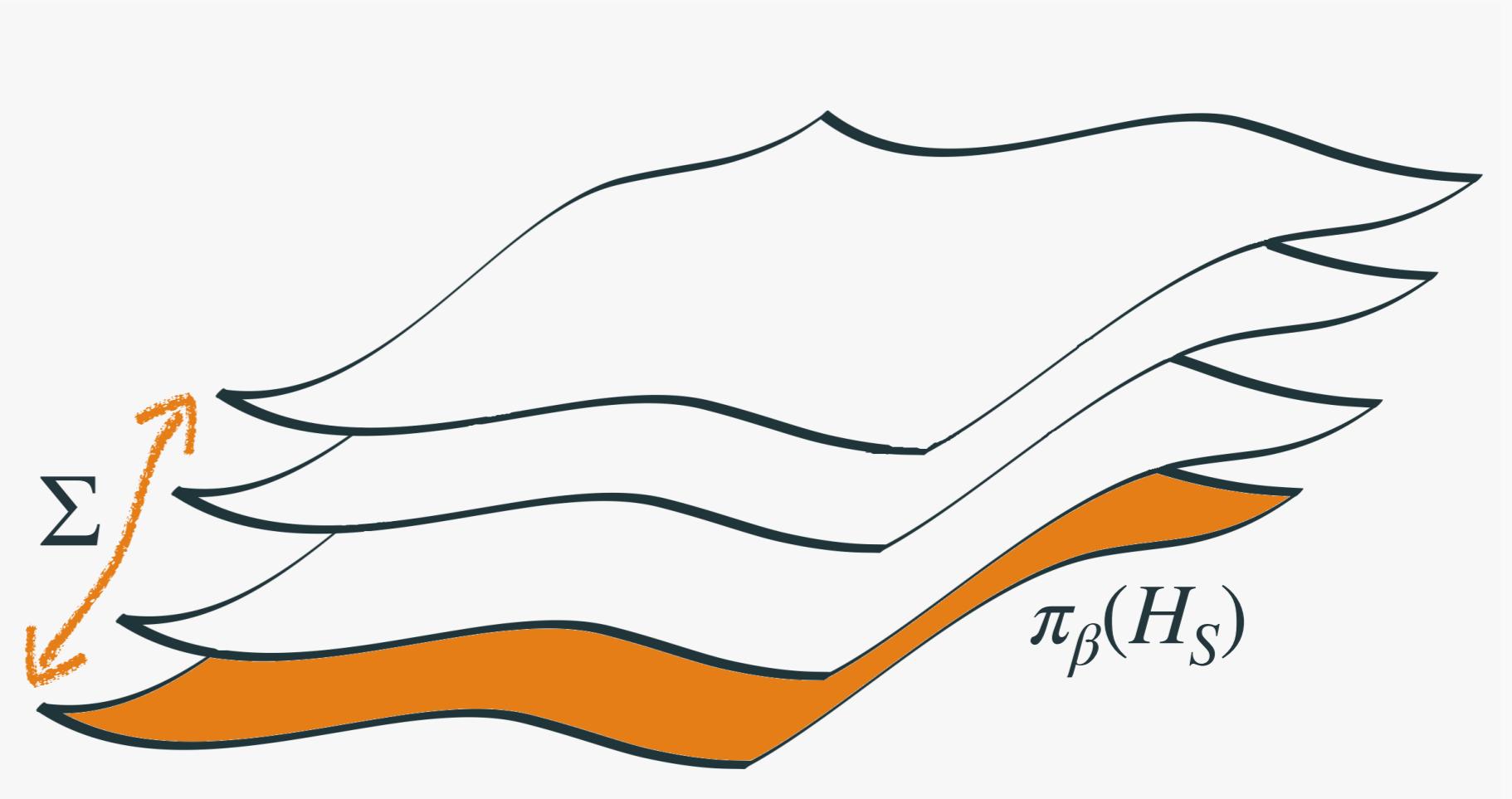
$$\begin{aligned}\beta Q &= \Delta S - \Sigma \\ w &= \Delta F + w_{diss}\end{aligned}\quad \boxed{\text{First Law}} \quad \Sigma = \beta w_{diss}$$

Thermodynamic space



$$H(t) = H_S(t) + gV + H_B$$

(H_S, ρ_S)



$$\rho_S \sim \pi_\beta(H_S)$$

$$\Sigma \sim x^i (\partial_i \partial_j \Sigma) x^j$$

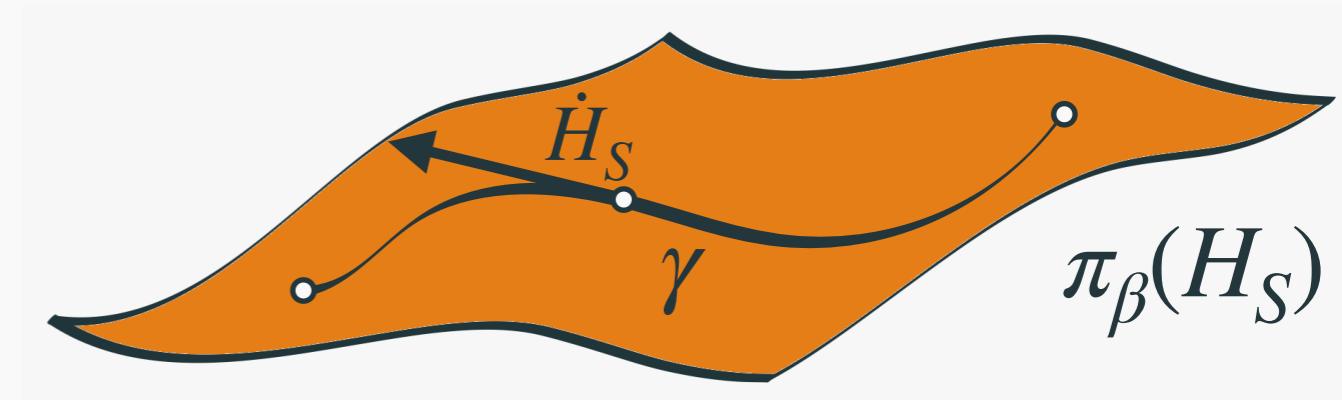
Close to equilibrium

Average dissipation¹



$$\dot{\rho}_S = \mathcal{L}_t[\rho_S]$$

$$T \gg \tau$$



$$\langle w_{diss} \rangle = - \frac{1}{T} \int_{\gamma} \int_0^1 dy \text{cov}_t^y(\dot{H}_t, (\mathcal{L}_t^+)^{\dagger}[\dot{H}_t])$$

$$\text{cov}_t^y(A, B) = \text{Tr}[\pi_t^{1-y} A \pi_t^y B] - \text{Tr}[\pi_t A] \text{Tr}[\pi_t B]$$

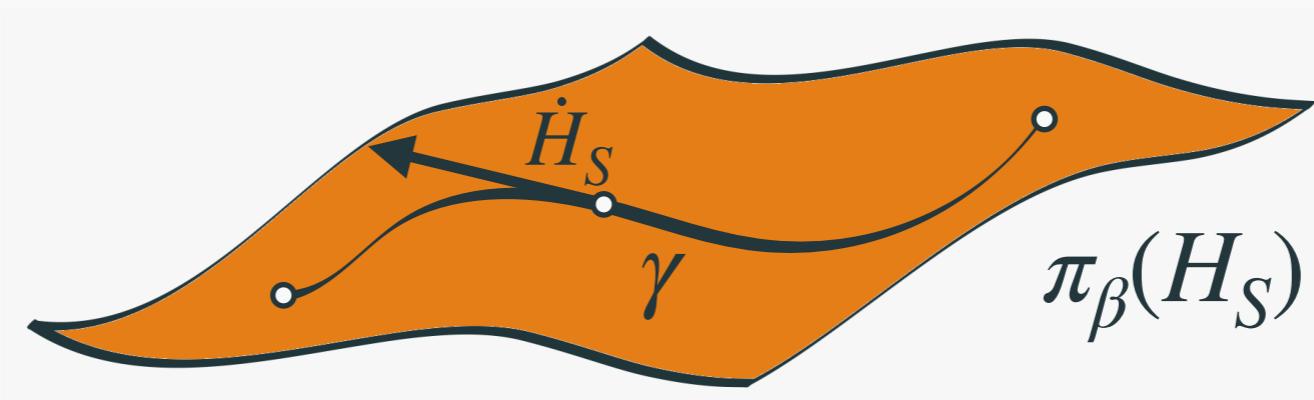


Thermodynamic metric¹

$$\int_{\gamma}^1 \int_0^1 dy \operatorname{cov}_t^y(\dot{H}_t, (\mathcal{L}_t^+)^{\dagger}[\dot{H}_t])$$

- Symmetric
- ≥ 0
- Smooth

Metric

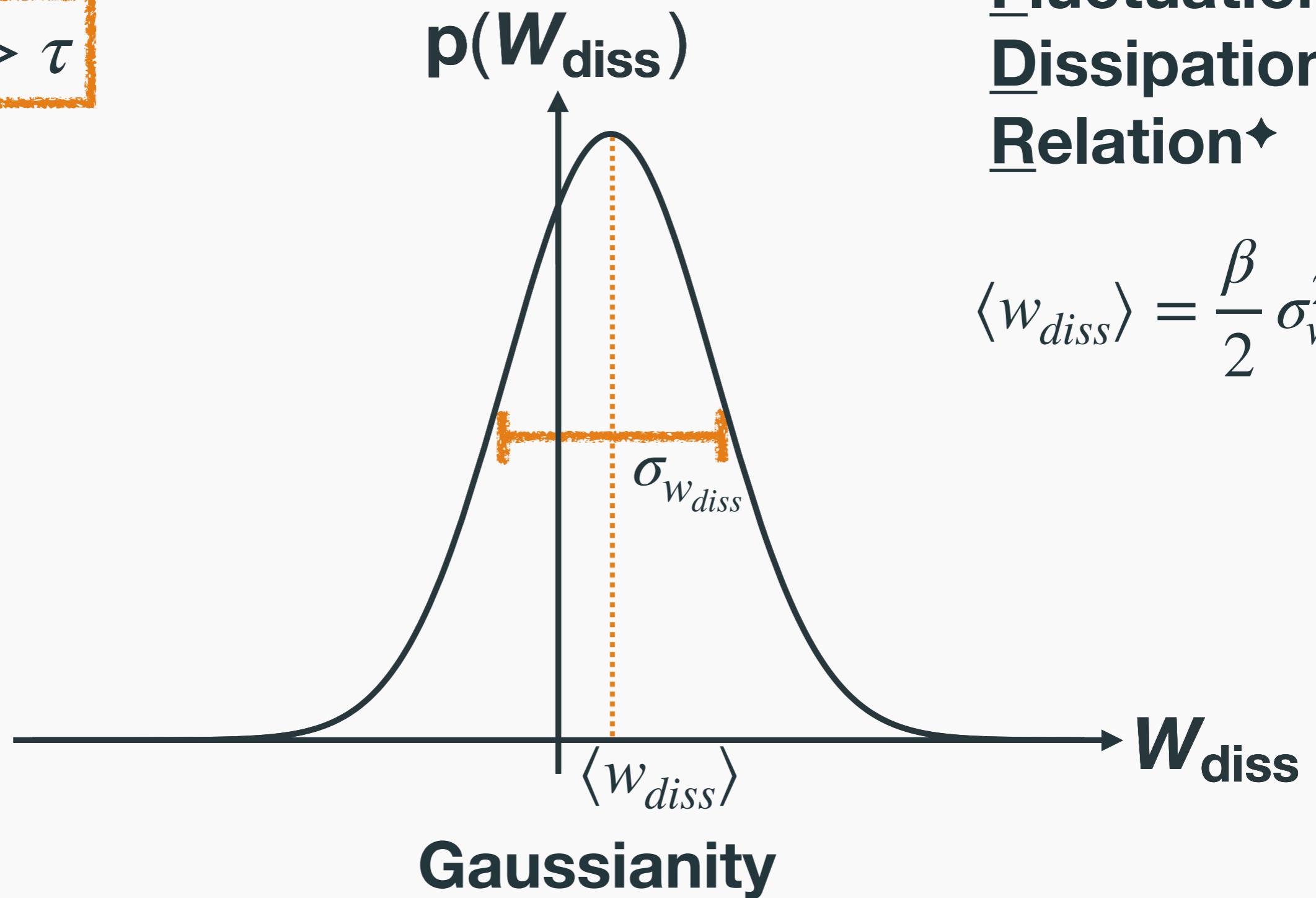


Geodesics minimise dissipation

Classical case



$T \gg \tau$



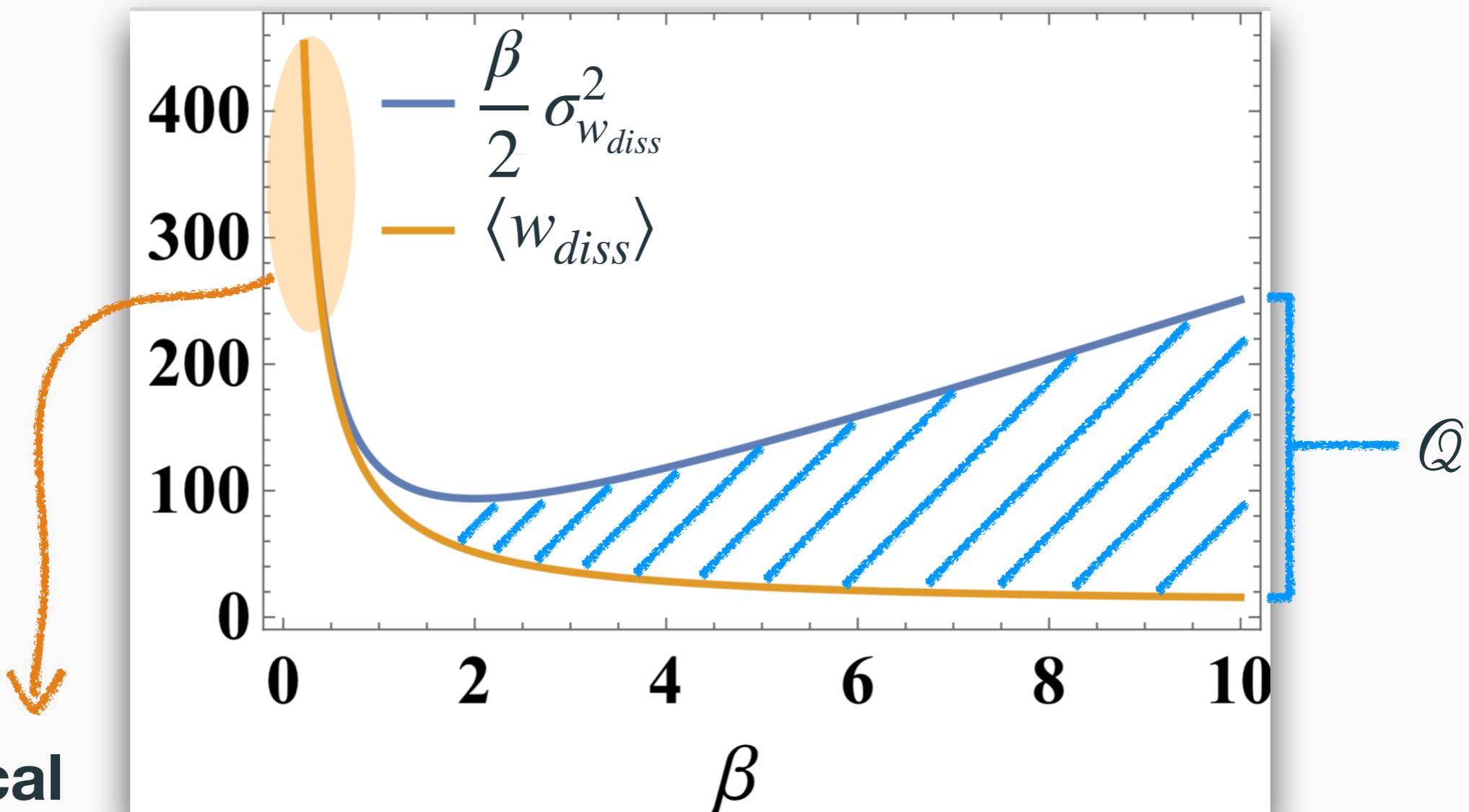
Fluctuations
Dissipation
Relation♦

$$\langle w_{\text{diss}} \rangle = \frac{\beta}{2} \sigma_{w_{\text{diss}}}^2$$

Quantum FDR²



$$\langle w_{diss} \rangle = \frac{\beta}{2} \sigma_{w_{diss}}^2 + Q$$



Cumulant Generating Function³



$$K(\lambda) = \log \int p(w_{diss}) e^{-\beta \lambda w_{diss}} dw_{diss}$$
$$= \sum_n \frac{\lambda^n}{n!} \kappa_n$$

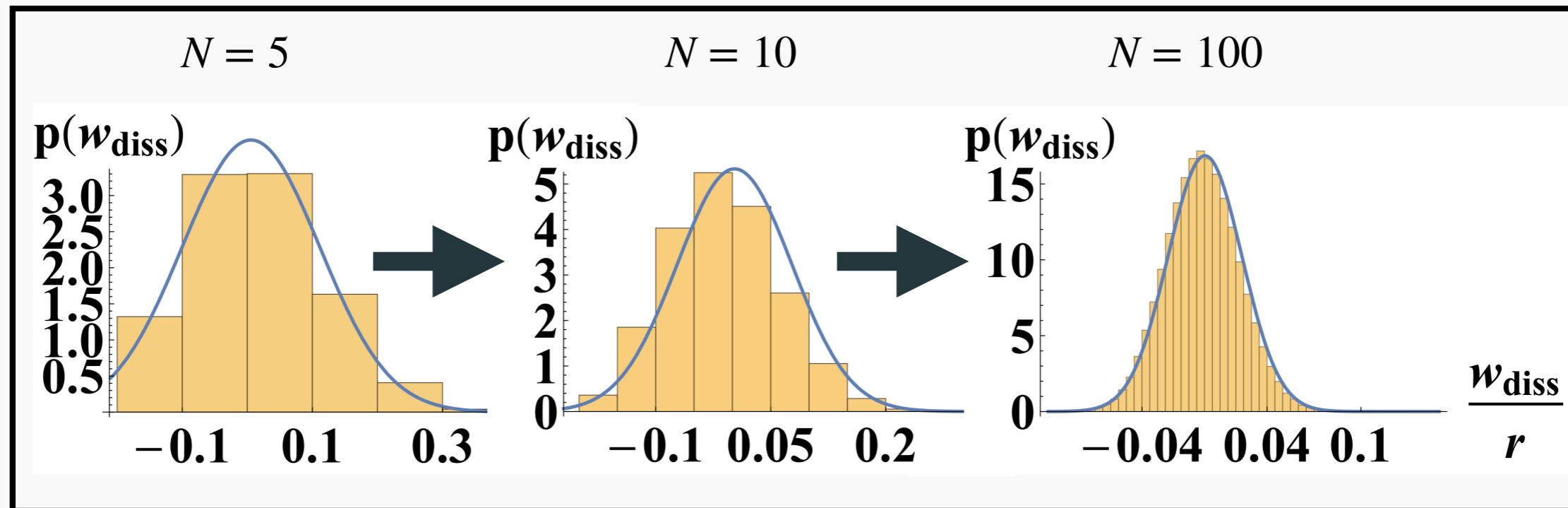
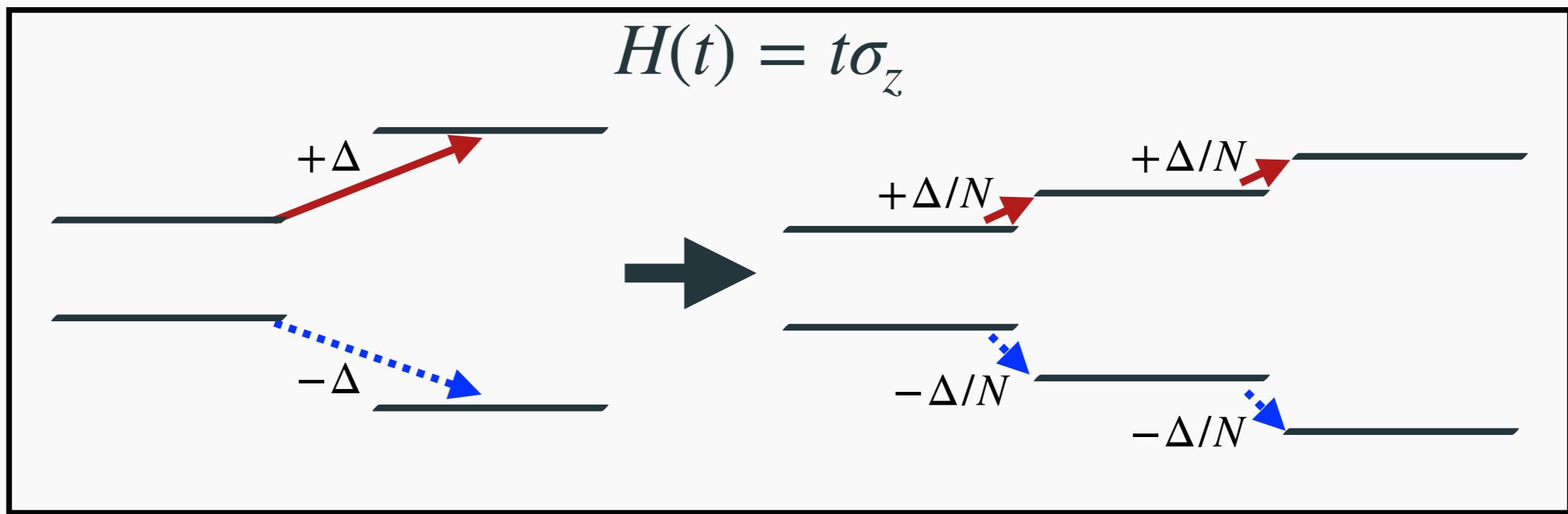
$$= (\lambda - 1) S_\lambda(\pi_T || \rho_T)$$

$T \gg \tau$

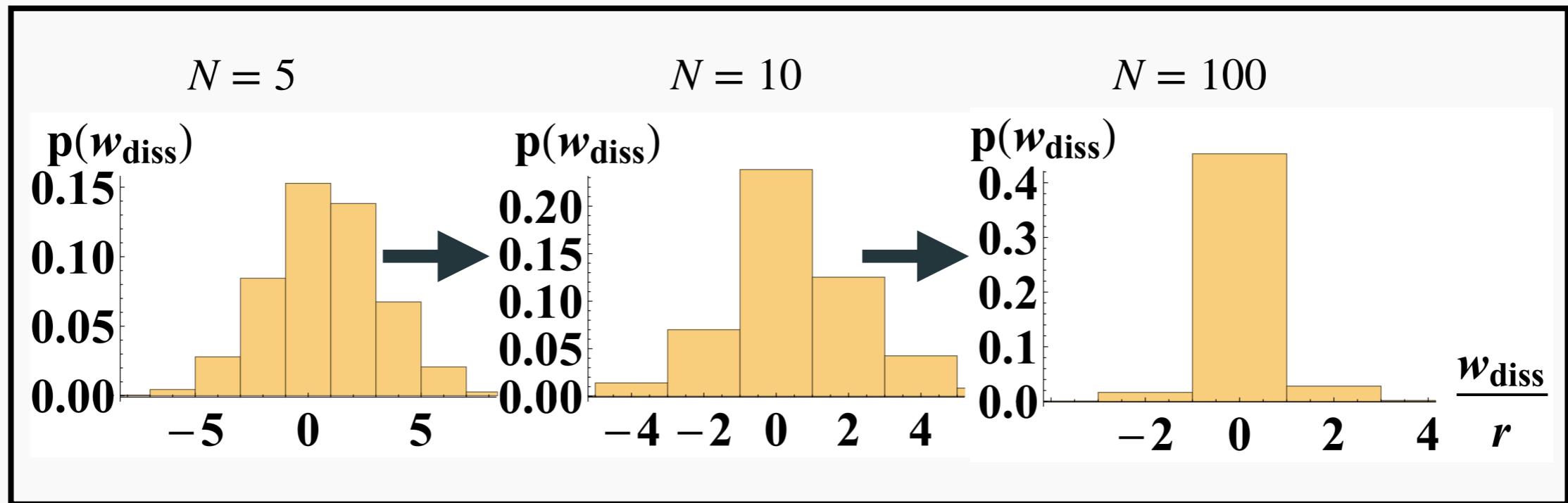
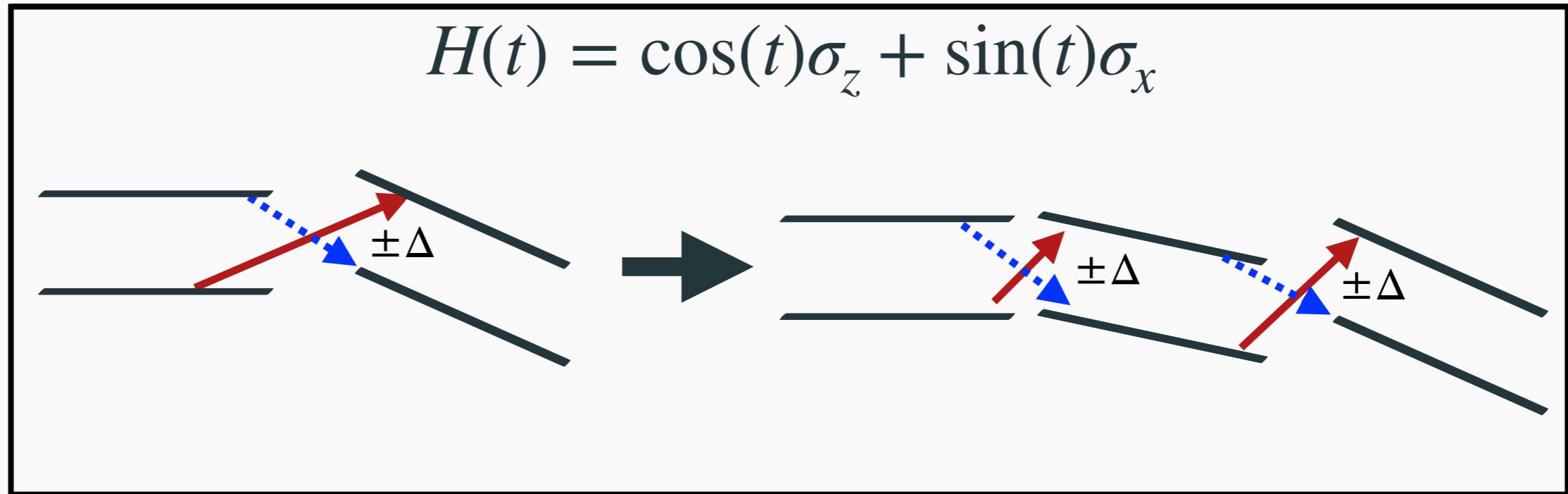
$$= \frac{\beta^2(\lambda^2 - \lambda)\tau}{2T} \int_\gamma \text{Var}_{\pi_t}[\dot{H}_t] + \frac{\beta^2\tau}{2T} \int_\gamma \int_0^\lambda dx \int_x^{1-x} dy I^y(\pi_t, \dot{H}_t)$$

Classical Quantum contribution

Qubit: diagonal driving

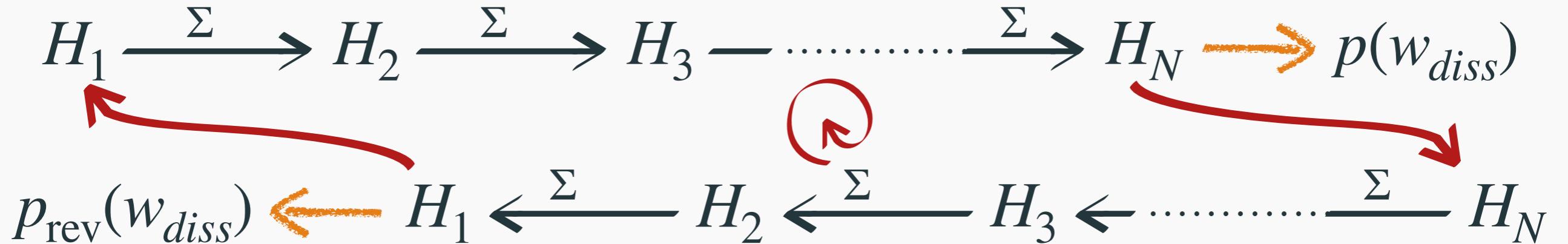


Qubit: coherent driving

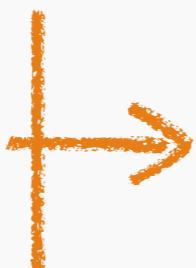


Simplifications

Time reversal



Crooks
relations

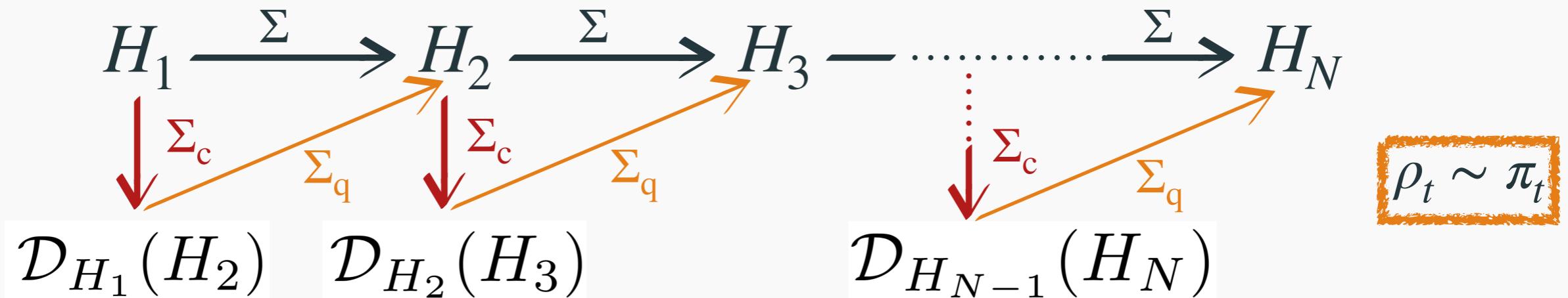


$$p(w_{diss}) = p_{\text{rev}}(-w_{diss}) e^{\beta w_{diss}}$$

$$T \gg \tau$$

$$p(w_{diss}) = p_{\text{rev}}(w_{diss}) \xrightarrow{\text{orange hand-drawn style arrow}} p(w_{diss}) = p(-w_{diss}) e^{\beta w_{diss}}$$

Channels of entropy production



Second laws

$S_\lambda(\rho || \pi)$ **decreases**

$S_\lambda(\rho || \mathcal{D}(\rho))$ **decreases**

Summary of the results



Quantum signatures

$$\frac{\beta}{2} \langle \sigma^2 \rangle = \langle W_{diss} \rangle + Q$$

Non gaussianity

Consequences of Slow driving

Time reversal symmetry

Decoupling of entropy production channels

Final remarks



- ❖ Quantum signatures survive the thermodynamic limit (FDR)
- ❖ Markovian dynamics define a family of metrics
- ❖ Extendible to strong coupling



**Witnessing
Non Markovianity**





Thanks for the attention!

Thermodynamic length: arXiv:1810.05583

Fluctuation dissipation relations: arXiv:1905.07328

Work statistics: arXiv:1911.04306